



Conformable Derivatives in Laplace Equation and Fractional Fourier Series Solution

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ABSTRACT

In this paper the solution of conformable Laplace equation, $\frac{\partial^\alpha u(x,y)}{\partial x^\alpha} + \frac{\partial^\alpha u(x,y)}{\partial y^\alpha} = 0$, where $1 < \alpha \leq 2$ has been deduced by using fractional fourier series and separation of variables method. For special cases $\alpha = 2$ (Laplace's equation), $\alpha = 1.9$, and $\alpha = 1.8$ conformable fractional fourier coefficients have been calculated. To calculate coefficients, integrals are of type "conformable fractional integral".

Keywords: Conformable fourier series, conformable derivative, conformable Laplace equation.

1 Introduction

History of "conformable" fractional derivative which is not really fractional derivative, based on limit definition, returns to five years ago (in 2014). However, about starting point of fractional derivative, it should be said that L'Hopital asked the question "What does derivative of order $\frac{1}{2}$ mean? namely $\frac{d^{\frac{1}{2}}f}{dx^{\frac{1}{2}}}$ " in 1695. Many researchers have been trying to generalize the concept of an ordinary derivative to fractional derivative. Most of definitions use an integral form such as Riemann - Liouville definition or Caputo definition. However, Khalil et

al. introduces the limit form for fractional derivative entitled "conformable fractional derivative" [1–5].

Definition 1.1. [1] Given a function $f : [0, \infty) \rightarrow \mathbb{R}$

$$T_\alpha(f)(t) = \lim_{\epsilon \rightarrow 0} \frac{f(t + \epsilon t^{1-\alpha}) - f(t)}{\epsilon}.$$

for all $t > 0, \alpha \in (0, 1)$. If f is α -differentiable in some $(0, a), a > 0$, and $\lim_{t \rightarrow 0^+} f^{(\alpha)}(t)$ exist, then define $f^{(\alpha)}(0) = \lim_{t \rightarrow 0^+} f^{(\alpha)}(t)$.

Khali et al. gives also via fractional cords a geometrical meaning to conformable fractional derivative [6]. In [7] Tarasov



proves that Conformable derivatives can not be considered as fractional derivatives of non-integer order. Thus, it is better that instead of “Conformable fractional derivative” we just call it “Conformable derivative”.

For a more general case $\alpha \in (n, n+1]$, and a natural number n we have the following definition:

Definition 1.2. [1] Let $\alpha \in (n, n+1]$, and f be an n -differentiable at t , where $t > 0$. Then the conformable fractional derivative of f of order α is defined as

$$T_\alpha(f)(t) = \lim_{\epsilon \rightarrow 0} \frac{f^{([\alpha]-1)}(t + \epsilon t^{([\alpha]-\alpha)}) - f^{([\alpha]-1)}(t)}{\epsilon}$$

where $[\alpha]$ is the smallest integer greater than or equal to α .

The conformable fractional derivative T_α satisfies all the properties in the following theorem.

Theorem 1.1. [1, 5] Let $\alpha \in (0, 1]$ and f, g be α -differentiable at a point $t > 0$. Then,

1. $T_\alpha[af + bg] = aT_\alpha(f) + bT_\alpha(g)$, for all $a, b \in \mathbb{R}$.
2. $T_\alpha(t^p) = pt^{p-\alpha}$ for all $p \in \mathbb{R}$.
3. $T_\alpha(\lambda) = 0$, for all constant functions, $f(t) = \lambda$.
4. $T_\alpha(fg) = fT_\alpha(g) + gT_\alpha(f)$.
5. $T_\alpha\left(\frac{f}{g}\right) = \frac{gT_\alpha(f) - fT_\alpha(g)}{g^2}$.
6. If, in addition, f is differentiable, then $T_\alpha(f)(t) = t^{1-\alpha} \frac{df}{dt}(t)$.

Definition 1.3. [1] The α -fractional integral of a function f starting from $a \geq 0$ is defined as

$$I_\alpha^a(f)(t) = I_\alpha^a(t^{\alpha-1}f) = \int_a^t \frac{f(x)}{x^{1-\alpha}} dx, \quad (1)$$

where the integral is the usual Riemann improper integral, and $\alpha \in (0, 1)$.

Following the definition of conformable derivative and conformable integral, conformable differential equations, conformable integral equations and differential-integral equations involving them are also discussed. The conformable heat fractional partial differential equation and conformable diffusion fractional partial differential equation have been solved in [8, 9]. Furthermore, in generalized format a generalization of classical Sturm-Liouville equation to conformable fractional Sturm-Liouville equation has been introduced by Bilender et al. [10]. The methods for solving these equations can be analytical, semi-analytical or numerical methods [11–17]. The Fourier series is one of the most important methods that are used in engineering and physical sciences to give analytic solution involving initial and boundary values. The method of fractional series is introduced by Abu Hamed and Khalil [2]. By using of two methods, including combining conformable fourier series and separation of variables, in this paper we solve conformable Laplace fractional partial differential equation.

2 Results and Discussion

In order to solve conformable fractional Laplace's equation by method of fractional Fourier series we first need to consider conformable Fourier coefficients which can be calculated by using definition of conformable fractional integral. Calculation of conformable Fourier series coefficients are somewhat more complex than ordinary Fourier series. This series can be written as follows:

$$S(f)(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} [a_n \cos(n \frac{x^\alpha}{\alpha}) + b_n \sin(n \frac{x^\alpha}{\alpha})], \tag{2}$$

where $a_0, a_n,$ and b_n are fractional Fourier coefficients.

$$a_0 = \frac{1}{p} I_\alpha^a(f)(t) \Big|_{a=0}^{t=p=(\alpha 2\pi)^{\frac{1}{\alpha}}} \tag{3}$$

$$= \frac{1}{(\alpha 2\pi)^{\frac{1}{\alpha}}} \int_0^{(\alpha 2\pi)^{\frac{1}{\alpha}}} f(t) \frac{dt}{t^{1-\alpha}} \tag{4}$$

$$a_n = \frac{2}{p} I_\alpha^a(f \cos n \frac{t^\alpha}{\alpha})(t) \Big|_{a=0}^{t=p=(\alpha 2\pi)^{\frac{1}{\alpha}}} \tag{5}$$

$$= \frac{2}{(\alpha 2\pi)^{\frac{1}{\alpha}}} \int_0^{(\alpha 2\pi)^{\frac{1}{\alpha}}} f(t) \cos(n \frac{t^\alpha}{\alpha}) \frac{dt}{t^{1-\alpha}}, \tag{6}$$

$n = 1, 2, 3, \dots \text{and}$

$$b_n = \frac{2}{p} I_\alpha^a(f \sin n \frac{t^\alpha}{\alpha})(t) \Big|_{a=0}^{t=p=(\alpha 2\pi)^{\frac{1}{\alpha}}} \tag{7}$$

$$= \frac{2}{(\alpha 2\pi)^{\frac{1}{\alpha}}} \int_0^{(\alpha 2\pi)^{\frac{1}{\alpha}}} f(t) \sin(n \frac{t^\alpha}{\alpha}) \frac{dt}{t^{1-\alpha}}, \tag{8}$$

$n = 1, 2, 3, \dots$

The definition of periodic function also differs from the ordinary case as follows:

Definition 2.1. [2] A function $f(t)$ is called α -periodical with period p if

$$f(t) = g(\varphi(t)) = g(\varphi(t) + \frac{p^\alpha}{\alpha}),$$

for all $t \in [0, \infty)$.

Example 2.1. The function $f(t) = \cos(\frac{t^\alpha}{\alpha})$ is α -periodic with period $(\alpha 2\pi)^{\frac{1}{\alpha}}$. For value of $\alpha = \frac{2}{3}$, the function with period $(\frac{4\pi}{3})^{\frac{3}{2}}$ has been plotted in Figure 1. .

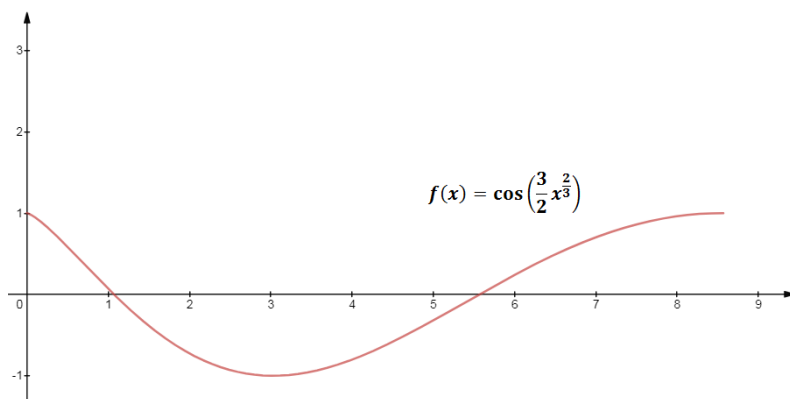


Figure 1: The function $f(x) = \cos(\frac{x^{\frac{2}{3}}}{\frac{2}{3}})$ is $\frac{2}{3}$ -periodic with period $(\frac{4\pi}{3})^{\frac{3}{2}}=8.573$.

Now we apply such series to solve the following conformable fractional Laplace equation after using separation of variables method:

$$\frac{\partial^\alpha u(x, y)}{\partial x^\alpha} + \frac{\partial^\beta u(x, y)}{\partial y^\beta} = 0, \quad (9)$$

where $1 < \alpha \leq 2$ and $1 < \beta \leq 2$; $\alpha = \beta$ along with the following boundary conditions:

$$\begin{aligned} u(0, y) &= u_1, \quad u(L, y) = u_2, \\ u(x, 0) &= u_3, \quad u(x, L) = u_4. \end{aligned} \quad (10)$$

We use the separation of variables method (SVM) so that $u(x, y) = X(x)Y(y)$. Putting it in the (9) gives

$$\frac{\partial^\alpha X(x)}{\partial x^\alpha} Y(y) + X(x) \frac{\partial^\beta Y(y)}{\partial y^\beta} = 0 \quad (11)$$

or

$$\frac{\frac{\partial^\alpha X(x)}{\partial x^\alpha}}{X(x)} + \frac{\frac{\partial^\beta Y(y)}{\partial y^\beta}}{Y(y)} = 0 \quad (12)$$

Since x and y are independent variables, then we should have

$$\frac{\frac{\partial^\alpha X(x)}{\partial x^\alpha}}{X(x)} = c, \quad \frac{\frac{\partial^\beta Y(y)}{\partial y^\beta}}{Y(y)} = -c,$$

where c is constant and should be determined.

So we get the solution

$$u(x, y) = e^{c(\frac{x^\alpha}{\alpha} - \frac{y^\alpha}{\alpha})}. \quad (13)$$

We can obtain the solution in terms of a double fourier series.

$$\begin{aligned} S(X)(x) &= \\ \frac{a_0}{2} &+ \sum_{n=1}^{\infty} [a_n \cos(n \frac{x^\alpha}{\alpha}) + b_n \sin(n \frac{x^\alpha}{\alpha})], \end{aligned} \quad (14)$$

$$\begin{aligned} S(Y)(y) &= \\ \frac{c_0}{2} &+ \sum_{m=1}^{\infty} [c_m \cos(m \frac{y^\alpha}{\alpha}) + d_m \sin(m \frac{y^\alpha}{\alpha})], \end{aligned} \quad (15)$$

$$a_0 = \frac{1}{(\alpha 2\pi)^{\frac{1}{\alpha}}} \int_0^{(\alpha 2\pi)^{\frac{1}{\alpha}}} e^{c \frac{x^\alpha}{\alpha}} \frac{dx}{x^{1-\alpha}} \quad (16)$$

$$\begin{aligned} a_n &= \frac{2}{(\alpha 2\pi)^{\frac{1}{\alpha}}} \int_0^{(\alpha 2\pi)^{\frac{1}{\alpha}}} e^{c \frac{x^\alpha}{\alpha}} \cos(n \frac{x^\alpha}{\alpha}) \frac{dx}{x^{1-\alpha}}, \\ n &= 1, 2, 3, \dots \end{aligned} \quad (17)$$

$$\begin{aligned} b_n &= \frac{2}{(\alpha 2\pi)^{\frac{1}{\alpha}}} \int_0^{(\alpha 2\pi)^{\frac{1}{\alpha}}} e^{c \frac{x^\alpha}{\alpha}} \sin(n \frac{x^\alpha}{\alpha}) \frac{dx}{x^{1-\alpha}}, \\ n &= 1, 2, 3, \dots \end{aligned} \quad (18)$$

And similarly we have

$$c_0 = \frac{1}{(\alpha 2\pi)^{\frac{1}{\alpha}}} \int_0^{(\alpha 2\pi)^{\frac{1}{\alpha}}} e^{-c \frac{y^\alpha}{\alpha}} \frac{dy}{y^{1-\alpha}} \quad (19)$$

$$\begin{aligned} c_m &= \frac{2}{(\alpha 2\pi)^{\frac{1}{\alpha}}} \int_0^{(\alpha 2\pi)^{\frac{1}{\alpha}}} e^{-c \frac{y^\alpha}{\alpha}} \cos(n \frac{y^\alpha}{\alpha}) \frac{dy}{y^{1-\alpha}}, \\ m &= 1, 2, 3, \dots \end{aligned} \quad (20)$$

$$\begin{aligned} d_n &= \frac{2}{(\alpha 2\pi)^{\frac{1}{\alpha}}} \int_0^{(\alpha 2\pi)^{\frac{1}{\alpha}}} e^{-c \frac{y^\alpha}{\alpha}} \sin(n \frac{y^\alpha}{\alpha}) \frac{dy}{y^{1-\alpha}}, \\ m &= 1, 2, 3, \dots \end{aligned} \quad (21)$$

We have calculated above coefficients for the special cases standard Laplace's equation (namely $\alpha = \beta = 2$), and conformable Laplace equation (Laplace-like equation) for values $\alpha = 1.9$ and $\alpha = 1.8$ in Appendix section.

3 Conclusion

This paper solves conformable Laplace equation, $\frac{\partial^\alpha u(x,y)}{\partial x^\alpha} + \frac{\partial^\alpha u(x,y)}{\partial y^\alpha} = 0$, by combination of fractional fourier series and separation of variables method. Fractional Fourier coefficients are of type conformable fractional integral that can be calculated in MAPLE software. To solve Conformable wave-like equation, $\frac{\partial^\alpha u(x,t)}{\partial x^\alpha} - \frac{1}{c^2} \frac{\partial^\alpha u(x,t)}{\partial t^\alpha} = 0$, conformable diffusion equation, $\frac{\partial^\alpha u(x,t)}{\partial x^\alpha} - \frac{1}{k} \frac{\partial^\beta u(x,t)}{\partial t^\beta} = 0$, in which $1 < \alpha \leq 2$ and $0 < \beta \leq 1$, one can follow this approach.

4 Competing Interests

The authors declare that there are no conflicts of interest regarding the publication of this paper.

5 How to Cite this Article:

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6 Appendix

1. Conformable fractional Fourier coefficients for $\alpha = 2$

$$\left[\begin{array}{l} > \frac{1}{(4\pi)^{\frac{1}{2}}} \int_0^{(4\pi)^{\frac{1}{2}}} x e^{c \frac{x^2}{2}} dx \\ > \frac{1}{(4\pi)^{\frac{1}{2}}} \int_0^{(4\pi)^{\frac{1}{2}}} x \cos\left(n \frac{x^2}{2}\right) e^{c \frac{x^2}{2}} dx \\ > \frac{1}{(4\pi)^{\frac{1}{2}}} \int_0^{(4\pi)^{\frac{1}{2}}} x \sin\left(n \frac{x^2}{2}\right) e^{c \frac{x^2}{2}} dx \end{array} \right. \quad \begin{array}{l} a_0 = \frac{1}{4} \frac{(-1 + e^{2\pi c}) \sqrt{4}}{c \sqrt{\pi}} \\ a_n = \frac{1}{4} \frac{(-c + 2e^{2\pi c} c \cos(\pi n)^2 - e^{2\pi c} c + 2e^{2\pi c} n \sin(\pi n) \cos(\pi n)) \sqrt{4}}{(c^2 + n^2) \sqrt{\pi}} \\ b_n = \frac{1}{4} \frac{(n + e^{2\pi c} n - 2e^{2\pi c} n \cos(\pi n)^2 + 2e^{2\pi c} c \sin(\pi n) \cos(\pi n)) \sqrt{4}}{(c^2 + n^2) \sqrt{\pi}} \end{array} \quad \begin{array}{l} (1) \\ (2) \\ (3) \end{array}$$

Figure 2: Fourier coefficients for solution of Laplace's equation $\frac{\partial^2 u(x,y)}{\partial x^2} + \frac{\partial^2 u(x,y)}{\partial y^2} = 0$.

2. Conformable fractional Fourier coefficients for $\alpha = 1.9$

$$\begin{aligned}
 &> \frac{1}{(3.8\pi)^{\frac{1}{1.9}}} \int_0^{(3.8\pi)^{\frac{1}{1.9}}} x^{0.9} e^{\frac{x^{1.9}}{1.9}} dx \\
 &\mathbf{a_0} = \frac{0.4952799403 \left(-1. + (e^{11.93805209 c})^{\frac{1052631579}{2000000000}} \right)}{c\pi^{0.5263157895}} \tag{4}
 \end{aligned}$$

$$\begin{aligned}
 &> \frac{1}{(3.8\pi)^{\frac{1}{1.9}}} \int_0^{(3.8\pi)^{\frac{1}{1.9}}} x^{0.9} \cos\left(n \frac{x^{1.9}}{1.9}\right) e^{\frac{x^{1.9}}{1.9}} dx \quad \mathbf{a_n} = \\
 &\frac{1}{(c^2 + n^2)\pi^{0.5263157895}} \left(0.4952799403 \left(-1. c + (e^{11.93805209 c})^{\frac{1052631579}{2000000000}} c \cos(6.283185310 n) \right) \right. \\
 &\quad \left. + (e^{11.93805209 c})^{\frac{1052631579}{2000000000}} n \sin(6.283185310 n) \right) \tag{5}
 \end{aligned}$$

$$\begin{aligned}
 &> \frac{1}{(3.8\pi)^{\frac{1}{1.9}}} \int_0^{(3.8\pi)^{\frac{1}{1.9}}} x^{0.9} \sin\left(n \frac{x^{1.9}}{1.9}\right) e^{\frac{x^{1.9}}{1.9}} dx \quad \mathbf{b_n} = \\
 &-\frac{1}{(c^2 + n^2)\pi^{0.5263157895}} \left(0.4952799403 \left(-1. n + (e^{11.93805209 c})^{\frac{1052631579}{2000000000}} n \cos(6.283185310 n) \right) \right. \\
 &\quad \left. - 1. (e^{11.93805209 c})^{\frac{1052631579}{2000000000}} c \sin(6.283185310 n) \right) \tag{6}
 \end{aligned}$$

Figure 3: Fourier coefficients for solution of Laplace-like equation $\frac{\partial^{1.9}u(x,y)}{\partial x^{1.9}} + \frac{\partial^{1.9}u(x,y)}{\partial y^{1.9}} = 0$.

3. Conformable fractional Fourier coefficients for $\alpha = 1.8$

$$\begin{aligned}
 &> \frac{1}{(3.6\pi)^{\frac{1}{1.8}}} \int_0^{(3.6\pi)^{\frac{1}{1.8}}} x^{0.8} e^{\frac{x^{1.8}}{1.8}} dx \\
 &\mathbf{a_0} = \frac{0.4908435123 \left(-1. + (e^{11.30973355 c})^{\frac{1388888889}{2500000000}} \right)}{c\pi^{0.5555555556}} \tag{7}
 \end{aligned}$$

$$\begin{aligned}
 &> \frac{1}{(3.6\pi)^{\frac{1}{1.8}}} \int_0^{(3.6\pi)^{\frac{1}{1.8}}} x^{0.8} \cos\left(n \frac{x^{1.8}}{1.8}\right) e^{\frac{x^{1.8}}{1.8}} dx \quad \mathbf{a_n} = \\
 &\frac{1}{(c^2 + n^2)\pi^{0.5555555556}} \left(0.4908435123 \left(-1. c + (e^{11.30973355 c})^{\frac{1388888889}{2500000000}} c \cos(6.283185307 n) \right) \right. \\
 &\quad \left. + (e^{11.30973355 c})^{\frac{1388888889}{2500000000}} n \sin(6.283185307 n) \right) \tag{8}
 \end{aligned}$$

$$\begin{aligned}
 &> \frac{1}{(3.6\pi)^{\frac{1}{1.8}}} \int_0^{(3.6\pi)^{\frac{1}{1.8}}} x^{0.8} \sin\left(n \frac{x^{1.8}}{1.8}\right) e^{\frac{x^{1.8}}{1.8}} dx \quad \mathbf{b_n} = \\
 &-\frac{1}{(c^2 + n^2)\pi^{0.5555555556}} \left(0.4908435123 \left(-1. n + (e^{11.30973355 c})^{\frac{1388888889}{2500000000}} n \cos(6.283185307 n) \right) \right. \\
 &\quad \left. - 1. (e^{11.30973355 c})^{\frac{1388888889}{2500000000}} c \sin(6.283185307 n) \right) \tag{9}
 \end{aligned}$$

Figure 4: Fourier coefficients for solution of Laplace-like equation $\frac{\partial^{1.8}u(x,y)}{\partial x^{1.8}} + \frac{\partial^{1.8}u(x,y)}{\partial y^{1.8}} = 0$.